





Numerical simulations of the collision of an inhomogeneous stellar wind and a relativistic pulsar wind in a binary system

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Outline

- 1. Introduction
- 2. Clump impact on the colliding wind region
- 3. Results
- 4. Discussion
- 5. Work in progress

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Introduction

We study the interaction of a **relativistic pulsar wind** and an **inhomogeneous stellar wind** in binary systems.



- The collision between the two winds creates a **shock structure**.
- Sources of stellar wind inhomogeneities:
 - Instabilities in the inner wind (Lucy & Solomon 1970).
 - Rotation, magnetic fields, or non-radial pulsations (Cranmer & Owocki 1996).
 - The truncation of the *Be* star equatorial **decretion disk** (Okazaki et al. 2011).
- The presence of **clumps** can **distort** the overall interaction structure (Bosch-Ramon 2013).
- This can affect the radiative output.

We perform numerical Relativistic Hydrodynamical simulations.

Introduction

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Clump impact on the colliding wind region I

The relativistic Euler equations can be written using the vector of conserved variables, U, and the flux vector, F as

$$\mathbf{U}_t + \mathbf{F}^{i}(\mathbf{U})_{x^{i}} = 0, \quad i = 1, 2, 3$$
 (1)

$$\mathbf{U} = \begin{bmatrix} \rho W \\ \rho h W^2 v^1 \\ \rho h W^2 v^2 \\ \rho h W^2 v^3 \\ \rho h W^2 c^2 - p - \rho W c^2 \end{bmatrix}, \quad \mathbf{F}^1 = \begin{bmatrix} \rho W v^1 \\ \rho h W^2 v^1 v^1 + p \\ \rho h W^2 v^2 v^1 \\ \rho h W^2 v^2 v^1 \\ \rho h W^2 v^2 v^1 \\ \rho h W^2 c^2 v^1 - \rho W c^2 v^1 \end{bmatrix}, \quad (2)$$

being the **conserved quantities** and v^i measured in the *laboratory frame*, and the **physical quantities** ρ and p measured in the *local rest frame* (e.g., Martí & Müller 2003).

The simulations were performed using a **finite-difference code** based on a high-resolution shock-capturing scheme that solves the multidimensional **equations of RHD in a conservation form** (Martí et al 1997).

- The code is parallelized using **OpenMP** (Perucho et al. 2005).
- The **intercell fluxes** are computed using Marquina's approach (see Donat & Marquina 1996).
- The **intercell physical quantities** are computed with a conservative monotonic parabolic reconstruction of the physical variables (**PPM**; see Woodward & Colella 1984; Martí & Müller 1996; and Mignone et al. 2005).
- The time integration is performed using a Runge-Kutta formula.

Clump impact on the colliding wind region III

We performed axisymmetric RHD simulations of the interaction of a relativistic pulsar wind and an inhomogeneous stellar wind (1 clump).



- Ideal gas with single particle species and $\hat{\gamma} = 1.444$ (cold protons and relativistic electrons).
- Grid size: $l_r = 2.4 \times 10^{12}$ cm, $l_z = 4.0 \times 10^{12}$ cm.
- The adopted resolution is 150×250 cells.
- The star is located at $(r_0, z_0) = (0, 4.8 \times 10^{12}) \text{ cm}.$
- The **pulsar** is placed at $(r_0, z_0) = (0, 0.4 \times 10^{12})$ cm.
- The star-pulsar separation is of $d_{\rm star-pulsar} = 4.4 \times 10^{12}$ cm.

Clump impact on the colliding wind region III

We performed axisymmetric RHD simulations of the interaction of a relativistic pulsar wind and an inhomogeneous stellar wind (1 clump).



- The lower and right **boundaries** are set to outflow. The left boundary is set to reflection.
- The **stellar wind** is injected as a boundary condition at the top of the grid.
- The **pulsar wind** is injected at a radius of 2.4×10^{11} cm (15 cells).
- The stellar mass-loss rate is $\dot{M} = 10^{-7} \ {\rm M_{\odot} \ yr^{-1}}.$
- The pulsar total luminosity is ${\it L}_{\rm sd}=10^{37}~{\rm erg~s^{-1}}~{\rm with}~{\rm Lorentz}~{\rm factor}~\Gamma=6.$
- The pulsar-to-stellar wind **thrust ratio** of $\eta \sim 0.2$.

Clump impact on the colliding wind region III

We performed axisymmetric RHD simulations of the interaction of a relativistic pulsar wind and an inhomogeneous stellar wind (1 clump).



- Set-up of the inhomogeneous stellar wind:
 - We obtained a quasi-steady state of the two-wind interaction considering an homogeneous stellar wind.
 - We introduced the stellar wind inhomogeneity characterized by a single clump centered at the z axis.
- We simulated 4 cases:

χ	$R_{ m clump}$	Description
10	$8\cdot10^{10}$ cm	light and small
10	$2\cdot 10^{11}$ cm	light and medium
10	$4\cdot 10^{11}$ cm	light and large
30	$8\cdot 10^{10}$ cm	dense and small

X. Paredes-Fortuny (UB)

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Results I



The results presented here have already been **published** in X. Paredes-Fortuny, V. Bosch-Ramon, M. Perucho, and M. Ribó (2015)

t = 58000 s



Clump: $\chi = 10$, $R_{clump} = 1 a = 8 \cdot 10^{10} cm$

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 670 s



Clump: $\chi = 10$, $R_{clump} = 1 a = 8 \cdot 10^{10} cm$

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X. Paredes-Fortuny (UB)



Clump: $\chi = 10$, $R_{clump} = 1 a = 8 \cdot 10^{10} cm$

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X. Paredes-Fortuny (UB)

t = 4000 s



Clump: $\chi = 10$, $R_{clump} = 1 a = 8 \cdot 10^{10} cm$

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 8700 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 1 a = 8 \cdot 10¹⁰ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 26000 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 1 a = 8 \cdot 10¹⁰ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 58000 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 2.5 a = 2 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 670 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 2.5 a = 2 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 2700 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 2.5 a = 2 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 5300 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 2.5 a = 2 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 11000 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 2.5 a = 2 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 18000 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 2.5 a = 2 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 58000 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 5 a = 4 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 670 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 5 a = 4 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 2000 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 5 a = 4 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 4000 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 5 a = 4 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 11000 s



Clump: $\chi = 10$, $R_{clump} = 5 a = 4 \cdot 10^{11} cm$

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

t = 15000 s



Clump: $\chi =$ 10, R $_{
m clump} =$ 5 a = 4 \cdot 10¹¹ cm

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 58000 s



Clump: $\chi = 30$, $R_{clump} = 1.0 a = 8 \cdot 10^{10} cm$

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 670 s



Clump: $\chi = 30$, $R_{clump} = 1.0 a = 8 \cdot 10^{10} cm$

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 2000 s



Clump: $\chi = 30$, $R_{clump} = 1.0 a = 8 \cdot 10^{10} cm$

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 4000 s



Clump: $\chi = 30$, $R_{clump} = 1.0 a = 8 \cdot 10^{10} cm$

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

X. Paredes-Fortuny (UB)

t = 8700 s



Clump: $\chi = 30$, $R_{clump} = 1.0 a = 8 \cdot 10^{10} cm$

Code units: $\rho_0 = 22.5 \times 10^{-22} \text{ g cm}^{-3}$, $a = 8 \times 10^{10} \text{ cm} \mid d_{\text{star-pulsar}} = 4.4 \times 10^{12} \text{ cm}$.

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General remarks:

- The **steady state** is very **sensitive** to the initial set-up parameters (grid size and density contrast / pulsar wind Lorentz Factor).
- The arrival of **clumps** can have very **strong impact** on the whole interaction structure (overcoming any possible numerical perturbation).
- The **clumps** trigger **Rayleigh-Taylor/Richtmyer-Meshkov** and **Kelvin-Helmholtz** instabilities leading to **quick changes** of the shocked pulsar-wind region (see also Bosch-Ramon et al. 2015).

Clump effects on the global structure and radiation:

- **Modest inhomogeneity degree**: non-negligible variations of the interaction structure and enhance the instability growth.
- **High inhomogeneity degree**: strong variations in the size of the two-wind interaction structure.
- Both generate quick and global variations in the shocked pulsar wind:
 - This affects the **location** of the pulsar wind **termination shock**, and therefore **reducing the emitter size** and increasing the magnetic field density (Bosch-Ramon 2013).
 - The relativistic **flow variations** on small spatial and temporal scales downstream of the pulsar wind shock **would lead to a complex radiative pattern in time and direction** caused by **Doppler boosting** (Khangulyan at al. 2014).

Discussion III

For the **PSR B1259–63 flare** observed by Fermi \sim 30 days after periaston passage (Abdo et al. 2011):

• The impact of a piece of disk on the two-wind interaction structure might have led to efficient Compton scattering by GeV electrons on local X-ray photons (Dubus & Cerutti 2013; Khangulyan et al. 2012 for IR photons), as a result of the strong enhancement of the synchrotron photon density.

For the **short flares of scales** of second to hours found in **X-ray** lightcurves of **LS 5039** and **LS I +61 303** (e.g., Bosch-Ramon et al. 2005; Paredes et al. 2007; Smith et al. 2009; Li et al. 2011):

• The impact of **clumps** might explain this short X-ray variability (Bosch-Ramon 2013).

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Work in progress — Radiative output



- V.M. de la Cita, V. Bosch-Ramon, X. Paredes-Fortuny, D. Khangulyan, M. Perucho, and M. Ribó in prep.
- Tomorrow's talk by Moreno de la Cita, Víctor at 9:50 on the same radiative code but applied to star-jet interactions in AGN.

Work in progress

Backup slides

Backup slides

Table 1: Wind velocity v, specific interal energy ϵ , and density ρ at a distance $r = 8 \times 10^{10}$ cm with respect to the star and pulsar centres located at (r_0, z_0) .

Parameter	Stellar wind	Pulsar wind
V	$3 imes 10^8~{ m cm~s^{-1}}$	$2.94 imes10^{10}~{ m cm~s^{-1}}$
ϵ	$1.8 imes 10^{15} \ { m erg g^{-1}}$	$9 imes 10^{19}~{ m erg~g^{-1}}$
ho	$2.68 imes 10^{-13} \mathrm{~g~cm^{-3}}$	$1.99 imes 10^{-19} { m g cm^{-3}}$
(r_0, z_0)	$(0, 4.8 imes 10^{12} { m cm})$	$(0, 4 imes 10^{11} ext{ cm})$

• Typical physical values in gamma-ray binaries:

 $egin{aligned} R_* &\sim 10 \, R_\odot, \ d_{
m star-pulsar} &\sim 10^{13} \ {
m cm}, \ R_{
m clump} &= &\sim 0.2 d_{
m star-pulsar} \end{aligned}$

High resolution simulations





Figure 2: Resolution 2 times higher. Steady state not reached.

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Work in progress

Results — light and <u>medium</u> size clump



Clump: $\chi=$ 10, $\mathsf{R}_{\rm clump}=$ 2.5 a = 2 $\cdot\,10^{11}$ cm

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